

Study of convective heat transfer in flat plate solar collectors

YEDILKHAN AMIRGALIYEV^{1,2}, MURAT KUNELBAYEV¹, ALIYA KALIZHANOVA^{1,2},
BEIBUT AMIRGALIYEV¹, AINUR KOZBAKOVA^{1,2}, OMIRLAN AUELBEKOV¹,
NAZBEK KATAEV¹

¹Institute Information and Computational Technologies CS MES RK, KAZAKHSTAN

²Al-Farabi Kazakh National University, KAZAKHSTAN

amir_ed@mail.ru, murat7508@yandex.kz, kalizhanova_aliya@mail.ru, ainur79@mail.ru,
omirlan@mail.ru, k_nazbek@mail.ru

Abstract: - The paper herein considers the study of convective heat transfer in flat plate solar collectors, as it is seen from the analysis on research of the heat transfer by a circular and flat tubes upon conforming the forced and free convections, placed vertically or horizontally with various liquid flow directions. There have been obtained Nusselt criterion dependencies in circular and flat pipes, which shows that corresponding equations allow defining the heat transfer intensity for all fluids with appropriate accuracy. As well there has been studied Prandtl criterion for fluids, which plays an important role in points spacing as regard to the curve. In each case the lines are with a tip 1/3, as in the research herein only the laminar conditions have been considered.

Key-Words: - Flat plate solar collector, convection, convective heat transfer, solar radiation

1 Introduction

To control and design many thermal systems it is necessary to specify heat transfer coefficients between a solid body and liquid [1], for instance, for solar collectors. Indeed, such thermal systems are always subject to temporal changes of some thermal parameters (solar thermal flow, liquid or solid substance temperature).

Therefore the heat transfer coefficient in the solid fluid medium cannot be considered as constant. Here is being investigated the problem of conjugated heat transfer [2]. In this case the heat transmission effect in the solid wall can't be neglected. If conjugated facts are neglected the temperature or thermal flow are under boundary conditions. However, in practice both the temperature and thermal flow on the solid body are defined simultaneously, which are linked with dynamic and energy equations [3].

It is known, that upon convective heat transfer computation according to experimental data, the effects of such heat transfer phenomena might be neglected, corresponding Nusselt numbers, as a rule, will be underestimated [4-6].

Recently, numerical methods have been applied to several researches [7-10]. Effect of bound conductive-convective flows in channels has been studied too in [11,12], with regular changes of temperature boundary conditions (temperature or thermal flow) or flow velocity. In the studies thereof it has been shown that thermal diffusion characteristic times in the solid body and fluid are

the parameters dominating in the bound heat transfer. There were carried out several researches dealing with the problems of the transition process heat transfer factors in the flat plates for various boundary conditions. Rebay & al. [13] submitted measurements of convective heat transfer local coefficients between the uniformly heated plate and air flow. Rebay & al. [14] studied unsteady forced convective flow heat transfer along a flat plate with a sharp change of thermal flow density on the surface without constant pressure along the flow direction.

Forced convection from microstructure on the flat plate has been investigated by Rebay & al. [15]. In [15] there has been studied the boundary layer longitudinal conductivity.

Differential method has been used to decrease differential equations with partial derivatives for ordinary differential ones, which are solved computationally by means of computer technology. Mladin & al. [16] simulated unsteady thermal boundary layer, being developed along the plate's finite thickness upon temperature change in the lower plate surface ramp. To simulate transient heat transfer there were applied two mathematical approaches: integral method based on the methodology of Carman- Pohlhausen equations system's NavierStokes complete equation. Both methods were in the mode of stationary state, numerical methods showed that both techniques are in line for 5% at steady condition and 2,6% at transient state. In the present work there has been used purely numerical method for studying the

unsteady laminar convection on the flat plate with different thickness subjected to the surface temporal change.

2 Method of research

2.1 Pilot tests of solar-driven installation

The goal of the current experiment is heating medium temperature determination in the heat absorber's heat removing channel.

Testing device is installed at the area latitude $\Psi = 43^\circ$, and the solar collectors slope angles equal to 55° , 45° , 30° in the circannian cycle of Kazakhstan, Almaty city.

Solar electric installation consists of accumulator box, solar collector, electric heater, support frame.

Accumulation tank has a heat-insulated cylinder case with a connecting branch for hot water intake and heat sealed stub-pipes for connecting the solar collectors.

On the tank cover there is a connecting branch for air outlet upon the system filling up.

Between the tank body and a case there is thermal insulation. At the tank bottom there mounted the pillars for fixing to the support. In the tank's lower part there installed an electric heating unit (EHU). Solar collector consists of black polyethylene tubes register with honeycomb insulation, lined out from above and below with pods (hydraulic collector), at that, the tubes length constitutes no less than 2,5 meters. Collector specific character is in the fact, that boundary tubes are manufactured from metal, and they allow regulating the thin-wall tubes strain, whereby enhancing stiffness of the tubular solar collector in the solar radiation absorption plane and excluding «housing unit».

To carry out the test at the solar site there is mounted SU (SI)-500 with tubular solar collectors. Tests have been conducted in the following mode: accumulator box (БА) was filled in with water at 9 am through a liquid level, i.e. filling in is an automatic process. Throughout the day there is water withdrawal from the accumulator tank. In the end of the working day – 5pm the water from accumulator tank is discharged completely.

2.2 Technical characteristics of solar collector

To determine the installation technical and operational characteristics there has been conducted temperature sensor readings recording. As temperature sensors there were used chromel-copel thermocouples to measure the temperature in 6

installation's points. Thermocouples readings recording has been fulfilled with an automatic recorder KCP-4. In parallel with thermocouples in a number of points there were installed mercury-filled thermometers for check measurements.

Following values were continuously recorded during the tests:

- flow density of the total solar radiation on the E collector surface;
- outside air temperature t_3 ;
- wind speed in unperturbed flow v ;
- cold tap water temperature t_3
- water temperature at the solar collector input and output t_1, t_2 ;
- water temperature in an accumulator box in three sections along the box height at a distance of 0,1, 0,3 and 0,4 m from the bottom t_5, t_6, t_7 ;
- temperature of the water, released to a consumer t_8 ;
- water discharge in the flow circuit G.

Measurements of all above mentioned parameters have been conducted every hour dawn-to-dusk starting from 9 am. Proceeding from measurement results there have been determined average total solar radiation flux density values during the tests on the \bar{E} collector surface, outside air temperature \bar{t}_3 , wind speed \bar{v} , water temperature at the collectors system input \bar{t}_1 and output \bar{t}_2 . Those values were found through planimetrying the squares, limited with curves $x(\tau)$, where x - is any of the denoted above values in the time range of the test conducting. Further there was computed thermal performance of the solar installation:

$$Q = G * c_1 * (t_2 - t_1) * \tau + \Delta Q_1$$

where τ -testing time; ΔQ_1 -heat, accumulated in-line by the experiment completion, which, as a rule, has not exceeded 0,05 and η is the solar installation daily performance

$$\eta = \frac{Q}{\sum E \cdot \sum S}$$

where $\sum E = \bar{E} \cdot \tau$ is the total amount of the sum solar radiation incident during the time τ on 1 m² on the solar collector surface and $\sum S$ is an overall square of the solar collectors heat absorption areas.

Tests outcomes for SU (SC)-500 with a tubular solar collector are listed in the table 1. Meteorological parameters average values constituted in the tests: $\bar{E}=375-888$ Wt/m²; $\bar{t}_3=14-20^\circ\text{C}$; $\bar{v}=1,0-4,5$ m/sec, a water flow in

water-flow circuit changed within 73,1-157,5 kg/hour. Consequently there has been received wide range of performance values from 0,2 to 0,64. Results of tests under usage of $G=2-3$ kg/sec. is approximated with linear dependence

$$\eta=0,782-11,3(\Delta \bar{t} / \bar{E})$$

Solar collector experimental findings thereof and computational fluctuations cannot be explained with

Date	Heating time		Water flow rate	Solar intensity	Air temperature	Temperature at collector inlet	Temperature at outlet	Inlet temperature at the tank	Resultant temperature in the tank	Installation heat producing	Installation performance	Specific heat producing	Heat production
	startin	final											
	τ_H	τ_K	G	E	t_0	t_1	t_2	t_3	t_4	Q	η	Q_{vL}	
	Hours, minutes		kg/hour	Wt/m ²	°C	°C	°C	°C	°C	VACH/hour.	%	VAC/m ²	
22.07.15	11:30	16:30	121,5	827	27,3	23,6	34,1	16	36	7400	0,59	493	525 150
23.07.15	11:50	17:50	157,5	888	27,5	25	31,7	18	38	7400	0,46	411	300
24.07.15	10:05	16:05	120,9	822	26,4	25	33,3	18	37	7030	0,47	390	
27.07.15	11:05	17:45	135	833	25,8	25,4	34,1	17	41	8880	0,54	455	150
28.07.15	11:30	16:30	130,7	883	25,6	27,4	37,1	19	39	7400	0,55	493	150
30.07.15	9:55	16:55	109	880	26,6	31,9	41,5	18	50	8537	0,46	406	
22.08.15	9:25	16:25	47,9	868	24,3	26,7	39	17	47	5324	0,29	253	
23.07.15	9:10	17:10	91,3	809	24,4	33,9	41,8	16,5	54	7090	0,36	295	170
26.08.15	9:30	14:30	110,7	665	25,6	21,4	28,1	16	29	4343	0,43	289	
27.08.15	10:15	17:15	110,3	920	25,9	35,4	42,2	27	48	6090	0,325	290	
14.09.15	9:30	16:30	73,1	916	26,4	30,7	40,5	20	49	5820	0,49	277	
16.09.15	14:45	17:10	97,9	824	27,4	48,1	53	44,5	53	6221	0,314	254	
7.09.15.	9:45	16:45	76,4	517	26	30,1	43,8	22	45	8510	0,48	405	145

As an example the figures 1 and 2 present the installation test results. During the test (from 9.00 to 17.00) a consumer received 0,5m³ of water with temperature about 50°C. There is defined the water temperature distribution along the accumulation tank height. That data illustrates well the water stratification extent in the tank. Temperature difference in the upper and lower strata amounted at the midday to 10-15°C.

Heat loss high level is conditioned by the fact, that the system's connecting pipelines total is about 5m and due to poor heat insulation, i.e., instead of polyurethane insulation there was used foam insulation.

2.3 Results of water temperature change upon different water volume

experimental errors. It may be shown, that upon parameter calculation $\Delta \bar{t} / \bar{E}$ directly on the water temperature at the outlet from the collector, all experimental points, corresponding to the run-off $G=2-3$ kg/sec., inclusive, in actual fact are on the one straight line.

Table 1 – Results of solar-driven installation with a tubular collector

Figures 1-7 present the results of water temperature change upon different water volume. The most optimal is the tank volume of 70-100 l per one square meter. With volume increase the water temperature in the tank lowers, for example, for 250 l–55-57°C, 300 l–50-53°C, 400 l–46-48°C, 500 l–45-47°C

According to the known data, at the latitude less than 54° there applied 90% of simple single circuit solar installation without a stand-by with natural circulation circuit for hot water supply of seasonal heat consumers.

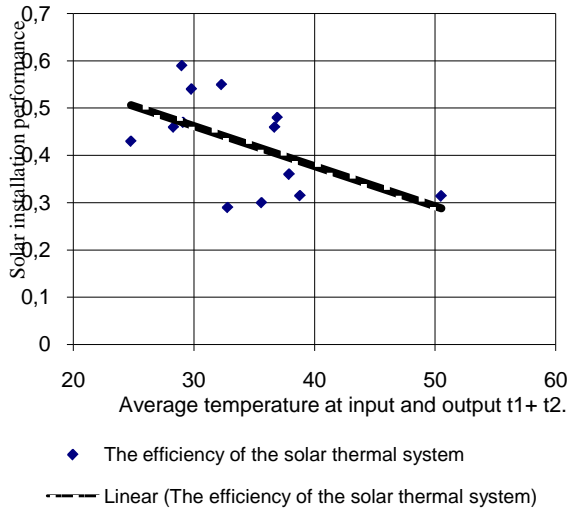


Fig. 1 Solar installation performance dependence on the average temperature at tubular collector input and output

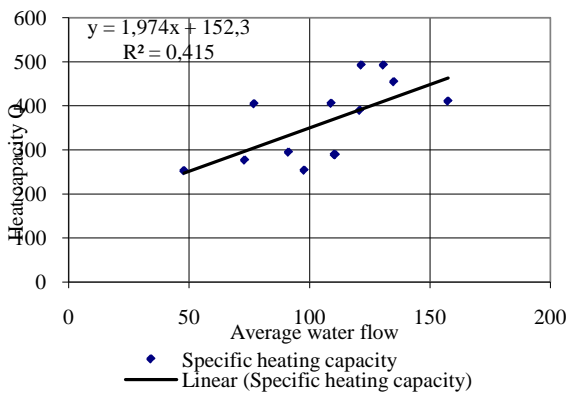


Fig. 2 Specific heat production dependence on the water flow in SI-500.

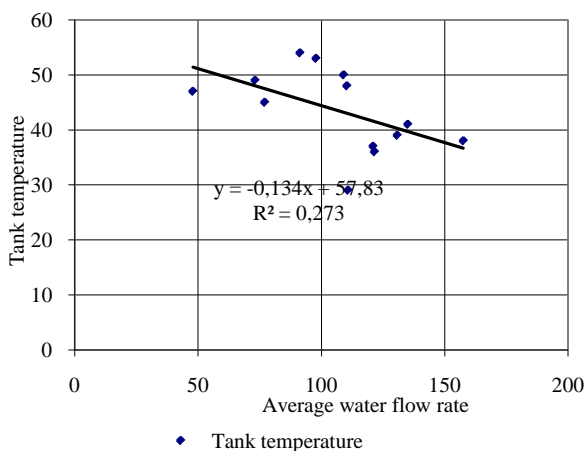


Fig.3 Dependence of heating temperature of water in the tank on water flow rate in the solar-driven installation's circulation loop

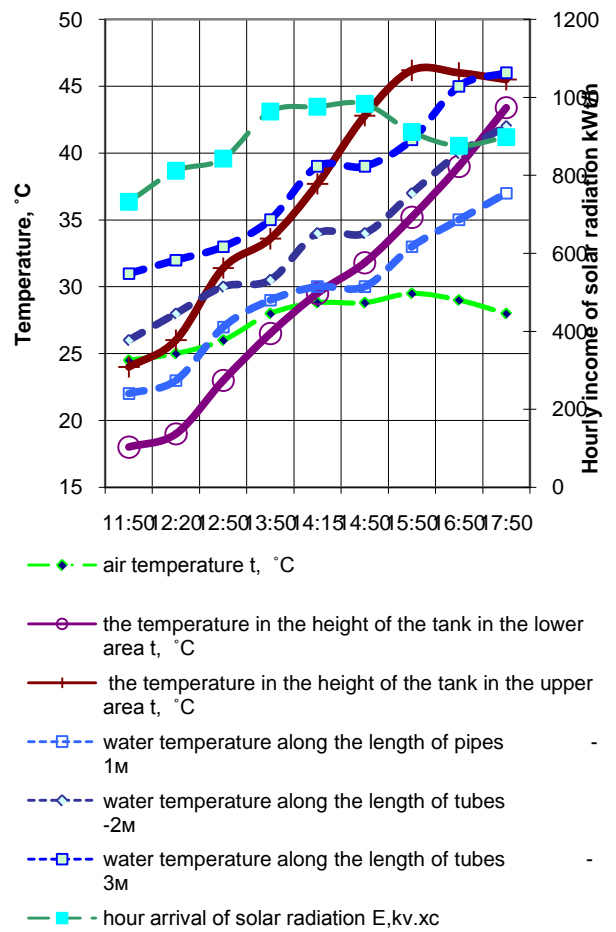


Fig. 4 Water heating dynamics in SU-500 with tubular collector with honeycomb transparent cover at water volume of 500 l

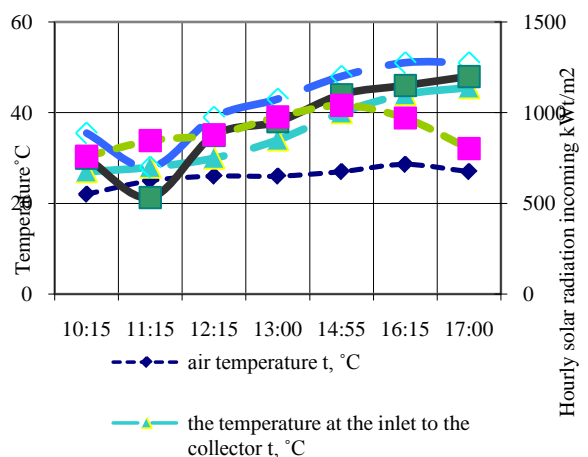


Fig.5 Water heating dynamics in SU-500 with tubular collector with honeycomb transparent cover at water volume of 400 l.

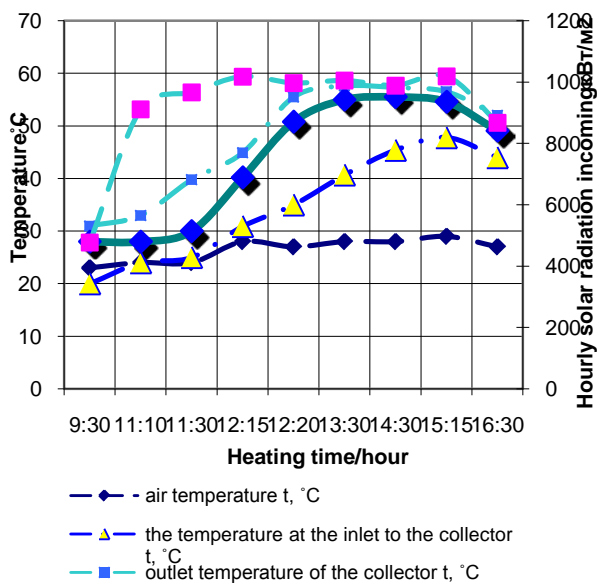


Fig. 6 Water heating dynamics in SU-500 with tubular collector with honeycomb transparent cover at water volume of 300 l.

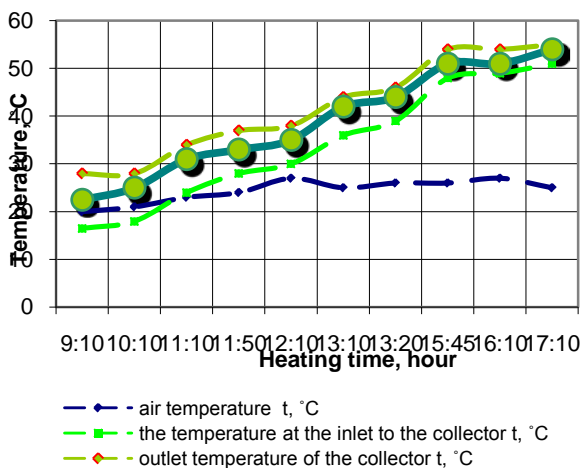


Fig. 7 Water heating dynamics in SU-500 with tubular collector with honeycomb transparent cover at water volume of 250 l.

In flat plate solar collectors heat absorption area and working liquid conditions are not favorable enough, particularly, when the heat receiver (boiler) is placed horizontally. Actually, the liquid motion speed is low here and it is heated from top, and, as a result, convective currents mixing the liquid do not develop. At the boiler's sloping position under known conditions there might happen the natural convection improving the heat transfer.

Task solution of heat transfer between radiant heat absorbing surface and power liquid flat heaters

upon applying the forced and natural convection (first type boundary conditions) presents sufficient difficulties. In case of applying the natural convection to the forced flow the velocity gradient on the wall will depend on the interrelated parameter part, defining both forced and free flow. In that case velocity gradient $A = \left(\frac{dW_A}{dy}\right)_{r,0}$ is not known in advance (for circular tube) $A = \rho \bar{W}/d$, for flat one $A = 6\bar{W}/h$, (A) can be defined from the motion equation.

Let us use the motion equation in a simplified form, omitting inertial terms, but taking into account the upward force. Supposing the liquid physical properties being constant and presenting the density in the form of linear temperature, we will obtain:

$$\frac{\partial^2 W_x}{\partial y^2} + \frac{g\beta(t_c - t_0)}{\nu} + \frac{1}{\mu} \frac{\partial P}{\partial x} = 0 \quad (1)$$

where β - volume coefficient; t_0 - flow temperature far from the wall, admitted as equal to the inlet temperature, as within the thermal initial section the temperature in the flow core changes ineffectually.

The equation for temperature distribution is given by:

$$\frac{t - t_c}{t_c - t_0} = \frac{\int_0^n e^{-\eta^3} d\eta}{\int_0^\infty e^{-\eta^3} d\eta} \quad (2)$$

where t_c - wall temperature, $\eta = \left(\frac{A}{gdx}\right)^{1/3} y$ - new independent variable. Denominator (2) represents the gamma function, the values of which are tabulated:

$$\int_0^\infty e^{-\eta^3} d\eta = \Gamma\left(\frac{3}{4}\right) = 0,8930 \quad (3)$$

Having substituted in (1) the value from (2) and upon integrating it we can define $A = \left(\frac{dW_x}{dy}\right)_{y=0}$. To carry out integration we approximate (2) and reduce the equation as follows:

$$\frac{t - t_c}{t_c - t_0} = 1 + b_1 e^{-\eta} + b_2 e^{-2\eta} + b_3 e^{-3\eta} \quad (4)$$

which meets boundary conditions and upon the constants appropriate selection conforms quite well with (1). Having inserted $t_c - t_0$ from (4) into (1), we find

$$\frac{\partial^2 W_x}{\partial y^2} \cdot \frac{g\beta(t_c - t_0)}{v} (1 + b_1 e^{-\eta} + b_2 e^{-2\eta} + b_3 e^{-3\eta}) - D \quad (5)$$

where $D = \frac{1}{\mu} \frac{\partial P}{\partial x} \approx const$

Boundary conditions are in the following form:
 $x = 0, t = t_0$

$$W_x = 6W \left(\frac{Y}{h} - \frac{Y^2}{h^2} \right), W_x = 8W \left(\frac{Y}{d} - \frac{Y^2}{d^2} \right) \quad (6)$$

The first equation is for flat, the second – for circular tube. Having executed integration and used boundary conditions we obtain the expression for non-dimensional velocity speed on the wall. For a flat tube

$$\varphi = \frac{Ah^2}{v} = 6Re + 0.43Gr \left(\frac{gx}{Pr\varphi h} \right)^{1/3} \quad (7)$$

$$Re = \frac{\bar{W}h}{v}; Gr = \frac{g\beta(t_c - t_0)h^3}{v^2}$$

Unfortunately, the expression thereof contains non evident φ . Therefore in (7) it was substituted with an approximate relationship:

$$\varphi = 6Re + 0.92 \frac{Gr^{3/4}}{Pr^{1/4}} \left(\frac{\alpha}{h} \right)^{1/4} \quad (8)$$

As you can see, the initial supposition about φ (or \bar{A}) consistency is not justified. However, φ changes in length relatively ineffectually ($\varphi \sim x^{1/4}$). Therefore, it will not be a great mistake if to use an average value in the section with the length:

$$\bar{\varphi} = \frac{\bar{A}h^2}{v} 6Re + 0.736 \frac{Gr^{3/4}}{Pr^{1/4}} \left(\frac{e}{h} \right)^{1/4} \quad (9)$$

Let's define now local heat-transfer coefficient, assigning it to the difference between the wall temperature and liquid temperature at the inlet to the heating area. Such definition technique (α) is convenient hereby, as at small values $\frac{1}{Pe} \frac{x}{d}$ the liquid average mass temperature changes weakly in length

$$\alpha_0 = \frac{g_c}{t_c - t_0} = - \frac{\lambda}{t_c - t_0} \left(\frac{dt}{dy} \right)_{y=0} \quad (10)$$

Having used the correlation obtained previously, we find:

$$\begin{aligned} \left(\frac{dt}{dy} \right)_{y=0} &= \left(\frac{A}{gax} \right)^{1/3} \left(\frac{dt}{d\eta} \right)_{\eta=0} \\ &= \left(\frac{A}{gax} \right)^{1/3} \left(\frac{t_c - t_0}{0.893} \right) \end{aligned} \quad (11)$$

consequently α_0 from (8) and (9)

$$\alpha_0 = \frac{\lambda}{0.893} \left(\frac{A}{9ax} \right)^{1/3} \quad (12)$$

Having substituted the value A from (10) in (11) we find the expression of Nusselt criterion for flow in a flat tube:

$$Nu = 1.467 \sqrt[3]{Pe \frac{h}{l} + 0.123 \left(GrPr \frac{h}{l} \right)^{3/4}} \quad (13)$$

Where $Nu = \frac{\alpha d}{\lambda}; Pe = \frac{\bar{W}d}{v}; Gr = \frac{g\beta(t_c - t_0)d^3}{v^2}$

In case of flow in a circular tube the similar analysis brings to an expression:

$$Nu = 1,615 \sqrt{Pe \frac{h}{l} + 0,092 \left(GrPr \frac{h}{l} \right)^{3/4}} \quad (14)$$

where $NU = \frac{\alpha d}{\lambda}; Pe = \frac{\bar{W}d}{v}; Gr = \frac{g\beta(t_c - t_0)d^3}{v^2}$

Heat transfer coefficient in (11) and (12) has been assigned to the initial temperature driving force $t_c - t_0$.

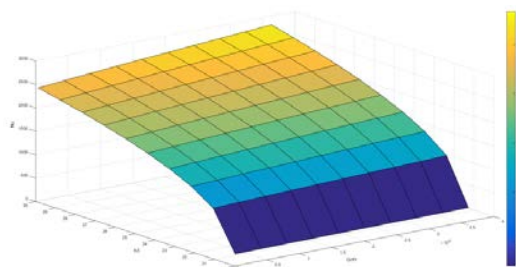


Fig. 8 Nusselt criterion for liquid flow in flat tube

Upon specifying the coefficient of heat transfer from the heating area to the heated water the thermal flow is defined according to the flow rate and temperature difference of the water being

heated at outlet and inlet. Water temperature is specified at outlet and inlet and the surface temperature through thermocouples readings average.

Experiments have been carried out at the heat transfer surface's obliquing angles $\phi=30^0, 35^0, 40^0$.

At that, the number $Re=150-500$, i.e., the experiments have been executed mainly at water flow laminar regime. As a determining temperature there has been accepted the water temperature and as a characteristic dimension – the channel equivalent diameter $d=4f/S$. Accepted in the model dimension d secures its geometric similarity to experimental flat solar collectors.

In the result, experimental data in the dependency form $Nu=f(Pld/l)$ and $Nu/Pr^{0.44}=f(Re)$ is given in the Figure 8.

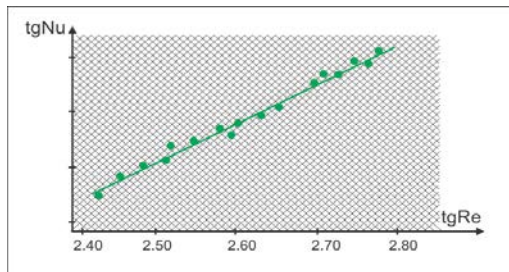


Fig. 9 Nusselt criterion dependence on Reynolds criterion for liquid flow in the flat solar collector

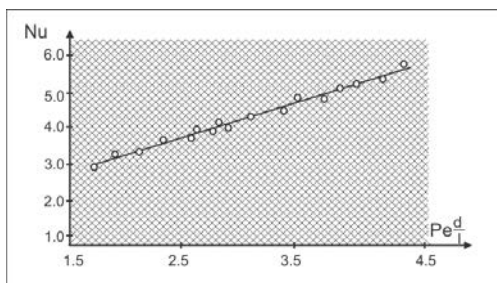


Fig. 10 Nusselt criterion dependence on Reynolds criterion for liquid flow in the flat solar collector

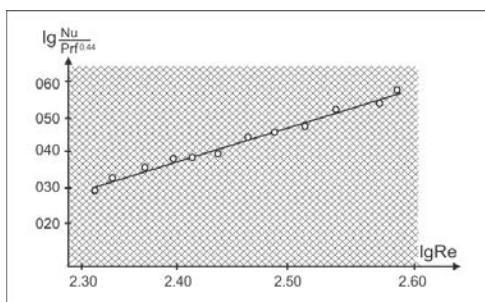


Fig. 11 Dependence of Nusselt-Prandtl criterion on Reynolds criterion for the liquid flow in the flat plate solar collector

Figures 9,10,11 show, that Prandtl number for the liquid plays an important role in the points spacing as regard to the curve. In all cases the lines are with a slope 1/3, as in the research herein we have considered only laminar mode.

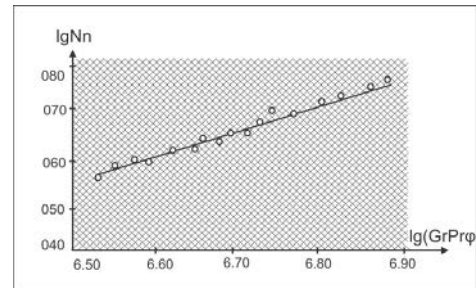


Fig. 12 Nusselt criterion dependence on Grashof - Prandtl criterion for liquid flow in flat- plate solar collector

Figure 12 demonstrates dependencies $lgNu=f[lg(GrPr)] \cdot \phi$ constructed with different obliquing angles. Proceeding from the considered plotted points for heat transfer medium, we can conclude, that in all cases of sloped position the free convection in the laminar flow at values of number $GrPr$ is higher than 10^5 . The Figure shows, that experiments outcomes are well described by the formula:

$$Nu=0,37(Pe \cdot d/l)^{0,35}(gr \cdot Pr \cdot \phi)^{0,18} \quad (15)$$

Obtained semiempirical formula allowing, in concept, processing and summarizing accumulated experimental data on the heat transfer upon convective heat transfer in the flat plate solar collectors as well gives a possibility to compare the results of theoretical investigations of heat transfer characteristics to experimental data.

Experimental data on heat transfer upon the mutual action of the forced and free convection can be broken down into three groups in compliance with the three main types of flow:

- a) flow in vertical tubes bottom-up at heating and top-bottom at cooling;
- б) flow in horizontal tubes at heating and cooling;
- в) flow in vertical tubes top-bottom at heating and bottom-up at cooling.

In the first case the forced and free convections at the wall match in direction, in the second case they are directed mutually perpendicular, in the third - opposite.

For the cases thereof there have been obtained corresponding equations, which allow defining the heat transfer intensity for all liquids with an

appropriate accuracy. Heat exchange conditions and experimental collected data of the researches thereof are given in the Table 2.

As it is seen from the above mentioned works analysis on studying the heat transfer with circular and flat tubes upon the forced and free convection conformity, ours and foreign researches have been conducted only with the tubes placed vertically or horizontally with different fluid flow directions.

Table 2

Mutual direction of forced and natural convection	Constants			Application boundaries		
	c	m	n	$Pe \frac{h}{l}$	$\frac{Gr \cdot h}{Pr}$	$\frac{h}{d}$
1	2	3	4	5	6	1
Conformity: flow in vertical pipes top-bottom at cooling and bottom-up at heating	0,3 5	0,3 0	0,8	110	$8 \cdot 10^5$ $4 \cdot 10^8$	20- 30
Mutually perpendicular direction: flow in horizontal tubes	0,3 5	0,6 0	0,1 0	200	$5 \cdot 10^5$ $1,3 \cdot 10^7$	60- 130
Opposite direction: flow in vertical tubes bottom-up at cooling and top-bottom at heating	0,21	0,80	0,00	300	$5 \cdot 10^5$ $1,33 \cdot 10^7$	60- 130

4 Conclusion

Obtained semiempiric formulae, which allow, in principle, processing and generalizing accumulated experimental data on heat transfer at convective heat transfer in flat plate solar collectors, as well they give a possibility to compare the theoretical researches outcomes of various heat transfer characteristics to experimental data.

This work is supported by grant from the Ministry of Education and Science of the Republic of Kazakhstan within the framework of the Project

«BR05236693 "Mathematical and computer models, hardware and software tools and experimental development on creation of network combined effective dual-circuit solar collectors with thermosiphon circulation and monitoring of their functioning».

References:

- [1] M. Rebay, A. Arfaoui, G. Mebarki, R. Ben Maad, and J. Padet, Improvement of the Pulsed photothermal technique for the measurement of the convective heat transfer coefficient, *Journal of Thermal Science*, Vol. 19, pp. 357-363, (2010)
- [2] P. R. Kanna, and M. K. Das, Conjugate forced convection heat transfer from a flat plate by laminar plane wall jet flow, *International Journal of Heat and Mass Transfer*, Vol. 48, pp. 2896-2910, (2005)
- [3] A. Pozzi, and R. Tognaccini, Conjugated heat transfer in unsteady channel flows, *International Journal of Heat and Mass Transfer*, Vol. 54, pp. 4019-4027, (2011)
- [4] Z. G. Liu, S. Q. Liang, and M. Takei, Experimental study on forced convective heat transfer characteristics in quartz microtube, *International Journal of Thermal Science*, Vol. 46, 139-148, (2007)
- [5] Celata, M. Cumo, V. Marconi, S.J. McPhail, and G. Zummo, Microtube liquid single-phase heat transfer in laminar flow, *International Journal of Heat and Mass Transfer*, Vol. 49, pp. 3538-3546, (2006)
- [6] G. L. Morini, Viscous heating in liquid flows in micro-channels, *International Journal of Heat and Mass Transfer*, Vol. 48, pp. 3637-3647, (2005)
- [7] C. Nonino, S. Savino, S. Del Giudice, and L. Mansutti, Conjugate forced convection and heat conduction in circular microchannels, *International Journal of Heat and Fluid Flow*, Vol. 30, pp. 823-830, (2009)
- [8] J. Y. Jang, and Y. W. Chiu, 3-D Transient conjugated heat transfer and fluid flow analysis for the cooling process of sintered bed, *Applied Thermal Engineering*, Vol. 29, pp. 2895-2903, (2009)
- [9] A. Barletta, E. Rossi di Schio, G. Comini, and P. D'Agaro, Conjugate forced convection heat transfer in a plane channel: Longitudinally periodic regime. *International Journal of Thermal Science*, Vol. 47, pp. 43-51, (2008)
- [10] A. Ates, S. Darıcı, and S. Bilir, Unsteady conjugated heat transfer in thick walled pipes involving two-dimensional wall and axial fluid conduction with uniform heat flux boundary

- condition, *International Journal of Heat and Mass Transfer*, Vol. 53, pp. 5058-5064, (2010)
- [11] B.A. Abu-Hijleh, M. A. Al-Nimr, and M. A. Hader, Thermal equilibrium in transient conjugated forced convection channel flow, *Numerical Heat Transfer A*, Vol. 43, pp. 327-339, (2003)
- [12] L. He, and M. L. G. Oldfield, Unsteady conjugate heat transfer modeling, *Journal of Turbomachinery*, Vol. 133, pp. doi:10.1115-1.4001245, (2011)
- [13] M. Rebay, M. Lachi, and J. Padet, Measurement of the convective heat coefficients by the pulsed method—influence of the boundary layer perturbation, *International Journal of Thermal Sciences*, Vol. 41, pp. 1161-1175, (2002)
- [14] M. Rebay, and J. Padet, Parametric study of unsteady forced convection with pressure gradient, *International Journal of Engineering Science*, Vol. 43, pp. 655-667, (2005)
- [15] M. Rebay, J. Padet, and S. Kakac, Forced convection from a microstructure on a flat plate, *Heat and Mass Transfer*, Vol. 43, pp. 309-317, (2007)
- [16] E. C. Mladin, D. Stanciu, and J. Padet, Transient thermal coupling for flows over a finite thickness plate exposed to a time-dependent temperature, *Proceedings of the Romanian academy*, Vol. 11, pp. 163-170, (2010).