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Numerical Investigation of Wall Shear Stress under Single-Phase and Two-Phase Flow in the Electrical Submersible Pump

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Abstract: - The effects of the wall shear stress on an Electrical Submersible Pump (ESP) was investigated in this paper. A CFD model in ANSYS Fluent was proposed to simulate actual single-phase and two-phase flow. The bottom hole pressure was minimized by utilizing the artificial lift methods. The flowing fluids in pumps and pipes causes shear stress on surface interacting. In multiphase flow application pump damages on head degradation as well as shear stress affects. The K- ω turbulence model and the multiphase Mixture approach with the sliding technique used to solve the Navier-Stokes equation. To study the effects of gas-liquid (air-water) flow on the ESP and the pump handle ability, the rotation speeds were varied while the other parameters were kept constant. The rotation speeds simulated were at 500, 900, 1500, 2000 and 2500 rpm meanwhile the water flow rate and gas flow rate were kept constant with 20 L/min and 1% fraction, respectively. The results obtained show that as the rotation speeds were observed, moreover the wall shear stress (WSS) increases. Although, the wall shear stress in both single-phase and two-phase flow were tend to increase as the blades length increased, however for the single-phase flow the WSS was found higher in all the simulated rotational speeds.

Keywords: Cavitation; CFD; Electrical Submersible Pump; Gas Pocket; Mixture Model; Pump Head; Turbulence Model; Two-Phase Flow; Wall Shear Stress.

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1 Introduction

Artificial lift mechanism is a process of minimizing Bottom Hole Pressure (BHP) in order to enhance the production rate. Viscous fluid engineers utilize concepts of BHP to produce positive-displacement downhole pump discovered in Progressive Cavity Pump (PCP) as well as downhole centrifugal pump. Electrical Submersible Pump (ESP) is a type of centrifugal pump and the second most utilized artificial lift mechanism; moreover, it is the highest amount in produced volume in two-phase flow. Applications involving Liquefied Natural Gas (LNG) face operational challenges under multiphase flow along with cavitation risk of pump component [1]. Oil and gas industry utilize artificial lift mechanism to boost produced fluid towards the surface during insufficient pressure. Similarly, an alternate artificial lift mechanism, known as gas lift is utilized for weight reduction of hydrostatic column through gas insertion in the tubing [2]. Pressure is the primary variable in artificial lift method, and is perpendicular to the surface; whereas, shear stress is parallel to target surface.

Centrifugal pumps transfer centrifugal forces to the working fluids, the fluids losses some of the transferred forces as shear stress between the blade wall and working fluid. Shear stress in continuum mechanics is defined as the stress applied parallel to the cross-section; however, in fluid mechanics, it refers to the effects of viscous flow on surface and between fluids in multiphase flow. Shear stress for fluids in motion is the rate of strain as shown in equation below:

$$\tau = \mu \frac{\partial \gamma}{\partial t} \tag{1}$$

where τ is shear stress, μ dynamic viscosity, and $\frac{\partial \gamma}{\partial t}$ strain rate.

Once fluid flows in a pipe, velocity of fluid at the fluid-pipe interface is equal to the pipe velocity (which is zero), and is known as no-slip condition. Under open wall flow, the region between boundary and height where the velocity of flow will equal the fluid velocity is known as boundary layer region. In Newtonian and laminar region, shear stress is proportional to strain rate, in which velocity gradient plays a dominant role as per the equation:

$$\tau(y) = \mu \frac{\partial u}{\partial y} \tag{2}$$

where $\frac{\partial u}{\partial y}$ is the velocity gradient.

However, study of wall shear stress requires the height at boundary assumed to be zero, so the integration is limited at the surface where the height is zero:

$$\tau(y=0) = \mu \frac{\partial u}{\partial y}\Big|_{y=0}$$
(3)

In Non-Newtonian flow, the viscosity is no longer a constant function with its axis expressed as:

$$\tau(\vec{u}) = \mu(\vec{u}) \,\nabla \,\vec{u} \tag{4}$$

This research investigates the effects of wall shear stress at the impeller blade walls of an ESP; furthermore, it observes the effects of gas bubbles upon the wall shear stress as well as the reduction in area of contact between the wall and the water.

2 Background Search

Gonzalez et al. (2002) [3] and Gonzalez & Santolaria (2006) [4] utilized CFD analysis to simulate the centrifugal pump. A transient model with sliding mesh technique was applied to solve the incompressible Navier-Stokes equation. The results obtained illustrated that the model implemented was able to predict the flow behavior inside the pump. Huang et al. (2010) simulated a three-dimensional single stage centrifugal pump by applying a transient model together with a sliding mesh technique and standard $k - \epsilon$ turbulence model. They investigated the interactions of flow between the impeller (rotor) and the diffuser (stator) inside the pump. The observed results illustrate an interaction pattern between rotor-stator that occurs due to pressure fluctuation and transient flow.

ANSYS Fluent offer various multiphase approaches to solve multiphase problems. The full Eulerian-Eulerian approach, the Volume Of Fluid (VOF) approach and Mixture approach. The difference between these approaches is applied momentum equation. Full Eulerian-Eulerian approach where n-terms of momentum equation are solved based on the number of phases is considered as the most advanced model. Meanwhile the VOF approach uses the fraction of the phases to solve one momentum equation [5]. The mixture model which is the model used in this paper is the summation of all momentum equations [6]. VOF and Mixture models provide accurate results in a time limit much less than the Eulerian-Eulerian model, due to smaller time step and finer mesh requirement for the Eulerian model.

Zhu et al. (2018) [7] used numerical simulation to simulate a three-stage pump in fluent to solve Reynolds Average Navier-Stokes equation (RANS). The Full Eulerian-Eulerian and VOF approaches with $k - \epsilon$ turbulence model were utilized to solve three different cases (case 1 λ_G = 0.35%, case 2 λ_G = 0.56% and case 3 $\lambda_G = 0.69\%$). The results implemented depict the full Eulerian model was more accurate than the VOF approach in terms of predicting the flow patterns and behavior of bubbles. Figure 1 represents the In-situ contour on second stage impeller for case 2 and case 3. In case 2 fig. 1-(a) an agglomerated flow pattern was observed due to coalescence of the bubbles at the impeller pressure side; however, increasing the gas fraction further as shown in fig. 2-(b), more coalescence of bubbles predicted which gas pocket was produced.



Fig.1: In-situ contour on second stage impeller, (a) Case 2 t= 7.956s, (b) Case 3 t = 15s [7]

Mizunuma and Nakajima (2007) [8], in the experimental study on shear stress distributions in a centrifugal Blood Pump, surface mounted hot film sensor was utilized to measure the wall shear stress and oil film visualization for shear distribution. From the hot film measurement, in the multiphase flow the shear stress was proportional to rotating speed, while in the boundary layer flow in the oil film method, the casing cover is a spiral flow from outside to inside. Shannak (2008) [9], performed experiments of twophase flow (air-water) on frictional pressure drop of vertical and horizontal as well as smooth and rough pipes, the friction pressure drop was observed to increase with an increase in the relative roughness of the pipe. In the study the parameters considered were the mass flux, the vapor quality, the density and viscosity of each phase, the pipe diameter, pipe orientation and the relative roughness of the pipe. In smooth and rough horizontal and vertical pipes, the frictional pressure drop increases with increasing mass flux, higher vapor quality and higher relative roughness. The influence of relative roughness is more significant especially at higher vapor quality and higher mass flux. Liu et al. (2019) [10] studied friction pressure drop in an inclined pipe at high velocity of two-phase flow (gas-water). The results obtained illustrate the pressure drop increases as the superficial velocities increase, the higher the velocity of the liquid phase the higher the contact with the wall, and the higher the gas flow velocity increases the higher the friction pressure drop. This led to develop a new model of friction pressure drop for different flow patterns.

Zubov et al. (2017) [11], stated that reliability of calculating friction pressure for two-phase flows close to atmospheric pressure in the field of twophase hydrodynamics still remains unsolved. The frictional pressure drop was calculated under low

pressure and compared with experimental results. Data from experiment included information on the flow mode boundaries, specific features of the individual structure, the profile of velocity, void fraction, droplet entrainment and deposition, on the distribution of phase interface surface in the flow, on the occurrence and characteristic of various types of instability and on the specific features of two-phase for small diameters. Under low reduced pressures, good results were obtained from using homogeneous model for quasi-homogeneous flow and the annular flow model for high values of void fraction. CFD modeling of flow profiles and interfacial phenomena in two-phase flow in pipes was studied by Ghorai and Nigam (2006) [12]. Two-phase (gas-liquid) flow in pipes was calculated using Eulerian-Eulerian model and commercial CFD package fluent. Parameters such as gas velocity, volume fraction of liquid, and interfacial roughness were considered in the study. Variation from 1.2 to 12.5 m/s were done on gas velocity and similarly, from 0.0066 m/s to 0.1 m/s in liquid. The velocity profile is symmetrical for low superficial gas velocity with a smooth interface; however, at higher velocities ranging from 4.7 m/s to 12.5 m/s the profile is no longer symmetric and hydro-dynamically interface became rough: moreover, the CFD analysis was validated with experiments on the concept of interfacial roughness for two-phase flow.

3 Methodology

3.1 Geometry

A replica of GDIWA A Series ESP is modeled as shown in the fig. 2 [13]. The geometry was designed accurately by using CAD program for CFD analysis in Fluent. The curves of impeller blades and diffuser outlet are the most effective parts and the most difficult on modeling, table 1 shows some details about impeller blades. The main components of the pump are impeller (rotor), Diffuser (stator), and Volute (fluid path) as show in the fig. 3.

Table 1: d	letails of	the imp	eller	blades
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Parameter	Symbol	Value
Blade height (mm)	Н	11
Number of blades	n	4
Impeller Inlet angle (°)	β_1	97
Impeller outlet angle (°)	β_2	44



Fig.2: Sectioning of GDIWA A Series ESP pump [13]



Fig.3: Assembled pump geometry for CFD analysis

3.2 Mesh Independency Test

Mesh independency test was performed to optimize the best mesh number that provides accurate results in a limited time frame. The performance curve of the pump was provided by the manufacture at 2850 rpm rotation speed [14]; moreover, 150 L/min flow rate and 2850 tetrahedron mesh type as shown in fig. 4 were selected as operating point for single-phase flow while the number of elements is varied by 1.5, 3, 5, 7 and 9 million elements. The results show that at 3 million elements the percentage of error dropped by 3% from 1.5 million elements while the simulating time is increased by 5 minutes; meanwhile, as the number of elements increased further the simulation time will increase and the reduction percentage is unfeasible due to trade-off with reduction in error. Fig. 5 shows the mesh independency test conducted for this research paper.





3.3 Setting up the Two-Phase flow Model

Two-Phase flow models are generally more complicated compared to single-phase flow model in defining parameters, since more boundary conditions are required to be defined in fluent; moreover, fluent offers various models for solving Navier-Stokes equation based on the study focused on. Table 2 represent more details of the models used and the boundary conditions specified for each component.

Table 2.	Details	of the	CFD	analysis	model
				~	

Model	Method
Turbulence model	K-ω SST
Multiphase model	Mixture model (implicit)
Body forces formulation	Implicit body forces
Interface	Sharp/Dispersed

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modeling	
Water, Air	Primary phase, Secondary
	phase
Boundary	Value
Conditions	
Velocity inlet	0.20938 <i>m/s</i> equivalent to 20
	L/min
Outlet boundary	Pressure outlet (gauge
-	pressure)
Rotational speed	500, 900, 1500, 2000, 2500
(impeller)	rpm
Solids boundary	Stationary relative to adjacent,
	no-slip condition
Water density	998 kg/m^3 , 99%
and fraction	
Air density and	$1.225 \ kg/m^3$, 1%
fraction	
Gravity	Y direction $g = -9.81 m/s^2$
Time step size	0.0001 s

This model consists of turbulence flow; therefore, it is crucial to focus the effects of turbulence on the wall. K- ω shear stress transport (SST) is a two equations eddy viscosity turbulence model that provides accurate results on adverse pressure gradient and near wall effects [15]. The standard K- ω equation is shown as the following:

$$\frac{\partial k}{\partial t} + U_j \frac{\partial k}{\partial x_j} = P_k - \beta^* k \omega + \frac{\partial}{\partial x_j} \left[(\nu + \sigma_k \nu_T) \frac{\partial k}{\partial x_j} \right]$$
(5)

where v_T is the kinematic eddy viscosity, and the transport equations of the *K*- ω SST are [16]:

$$\frac{\partial}{\partial t}(\rho k) + \frac{\partial}{\partial x_i}(\rho k u_i) + Y_k - S_k$$
$$= \frac{\partial}{\partial x_j} \left[\left(\mu + \frac{\mu_t}{\sigma_k} \right) \frac{\partial k}{\partial x_j} \right] + \tilde{G}_k \quad (6)$$

$$\frac{\partial}{\partial t}(\rho\omega) + \frac{\partial}{\partial x_i}(\rho\omega u_i) = \frac{\partial}{\partial x_j} \left[\left(\mu + \frac{\mu_t}{\sigma_\omega} \right) \frac{\partial \omega}{\partial x_j} \right] + G_\omega - Y_\omega + D_\omega + S_\omega$$
(7)

where

 \tilde{G}_k

Generation of turbulence kinetic energy due to mean velocity gradient G_{ω} Generation of ω $Y_k \& Y_{\omega}$ Dissipation of k and ω D_{ω} Cross-diffusion term $S_k \& S_{\omega}$ User defend source terms

The multiphase models are varied in the continuity equation, for the Mixture model the continuity equation form as:

$$\frac{\partial}{\partial t} \sum_{k=1}^{n} (\alpha_k \, \rho_k) + \, \nabla \cdot \sum_{k=1}^{n} (\alpha_k \rho_k u_k) = 0 \tag{8}$$

where α is the fraction and the subscript k is the phase. The general Navier-Stokes Equation can be written as the following:

$$\frac{d}{dt} \int_{V} \rho \phi \, dV \, \int_{\partial V} \rho \phi \left(\vec{u} - \vec{u_g} \right) . d\vec{A} = \int_{\partial V} \Gamma \, \nabla \phi . d\vec{A} + \int_{V} S_{\phi} \, dV \tag{9}$$

4 Results and Discussion

Figure 6 represents the results of air volume fraction at mid span of impeller blades, 20 L/min water flow rate and 1% air volume fraction. At low rotational speed of 500 rpm in fig. 6-(a) a large gas pocket is concentrated at the impeller inlet which choke the flow inside the pump with bubbles accumulating the blades pressure side to produce a gas pocket [17]. This leads to reduction in kinetic energy utilized to breakdown the bubbles concentration. As the rotational speed increases further as illustrated in fig. 6-(b), the large pocket begins to collapse to smaller bubbles with limited volume fraction at the blades pressure side. For higher rotation speeds 1500 rpm and 2000 rpm the gas pockets collapsed to small bubbles due to decrease in bubbles size as shown in fig. 6-(c) and (d); moreover, the bubbles moved towards the volute with higher fraction of 2%. Fig. 6-(e) illustrate the air volume fraction contour at maximum rotation speed 2500 rpm. It is clearly visible that large gas pocket and the concentration of bubbles at blades are collapsed totally due to high kinetic energy produced from the rotated impeller, and the high drag force applied by the liquid phase on the gas phase [18].



Fig.6: Air Volume Fraction Contours at mid span impeller blades

The results in this section focus on the effects of wall shear stress (WSS). WSS behavior on the pump performance plays a dominant role, as the rotation speed increases the friction pressure also increase [10]. Fig. 7 illustrates the Wall Shear Stress comparison of single-phase flow results at the impeller blade. The light blue color represents the 500-rpm rotation speed, pink is 900 rpm, black is 1500 rpm, red is 2000 rpm and dark blue is 2500 rpm. It can be observed that WSS is proportional to the rotation speed due to increases on the velocity gradient as per the equation below:

$$\tau = \mu \frac{\partial u}{\partial \nu} \tag{10}$$

The velocity gradient at the impeller blade increases as the length of the blade increases, which in turn results in the increase of wall shear stress. The values of WSS are fluctuating; however, the overall trend of the graph depicts an increase in WSS.



Fig.7: Wall Shear Stress of single-phase flow at different rpm's vs blade length

Two-phase flow wall shear stress has the same trend that of the single-phase at different rotational speed. Figure 8 below depicts the two-phase flow results of WSS vs the blade length. At different rotational speed from the plot of WSS vs blade length when the blade length beyond 85mm the wall shear stress drops rapidly due to the high velocity at tips as well as the flow separation and recirculation occurring at the blade tips [19]. Figure 9 shows the velocity vector at mid span of impeller blade of the 1500 rpm rotation speed and two-phase flow. The zoom in figure show that the velocity at the blade tip always recirculates due to the high velocity of the blade as well as the generated wake region [20].



Fig.8: Wall Shear Stress of two-phase flow at different rpm's vs blade length

Figure 10 illustrate the association between single and two-phase flow wall shear stresses at 1500 rpm. The blue color denotes the single-phase flow and red for the two-phase flow. Both trends increase the WSS values as the blade length increases; however, singlephase flow values are higher than that of two-phase flow, which is due to bubble's concentration at the impeller blade plus the viscosity effects between both flows [21]. The two-phase flow results show that air bubbles take place around the impeller blades, when the impeller rotates the bubble restrict the area of contact between the liquid (water) and the wall, which leads to reduction of wall shear stress. The detailed over view of this is illustrated in fig. 11, where a high concentration of air fraction around the blade root is more than the blade tip. Furthermore, in two-phase flow the viscosity is no more constant due to mixture viscosity [22].



Fig.9: Velocity vector at mid span of the blade, 1500 rpm and two-phase flow



Fig.10: Wall Shear Stress for 1500 rpm at different flow vs blade length



Fig.11: Detailed over view of bubble formation at 1500 rpm

5 Conclusion

The model simulated in this research was able to predict wall shear stress effects on the impeller blades for single-phase and two-phase flow. The K- ω SST Turbulence model and the mixture approach multiphase model with the sliding mesh technique

were utilized to visualize the flow patterns inside the pump. The rotation speed varied by 500, 900, 1500, 2000 and 2500 rpm while the water flow rate and gas fraction were kept constant with 20 L/min and 1%, respectively in all simulation. Finally, the results depict that the single-phase flow wall shear stress values are always higher than two-phase flow in all rotational speeds.

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- Abdulqader Hasan and Salman carried out Conceptualization, data curation, validation, and writing origin draft.
- Mohammed Alavi carried out methodology and writing origin draft.
- Dr. Sharul Sham Dol, Dr. Mohamed S. Gadala and Dr. Mohd Shiraz Aris carried out project supervision and project administration.

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